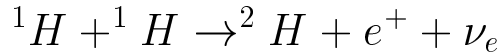


Nuclear reactions: proton-proton reaction

Nuclear reactions of interest generally involve several steps. In such a chain,

- The **rate** is controlled by the rate of the slowest reaction in the chain.
- The total energy release is given by the sum of the energies of all the individual steps.

The first, and rate determining step, in the proton-proton chain is,



where ${}^2_1\text{H}$ is a deuteron. This reaction is non resonant and slow because it involves a weak decay. The rate evaluates to,

$$r_{pp} = 1.15 \times 10^9 T_9^{-2/3} X^2 \rho^2 e^{-3.380/T_9^{1/3}} \text{ cm}^{-3} \text{ s}^{-1}$$

where T_9 is temperature in units of 10^9 K.

The temperature sensitivity is,

$$\nu_{pp} = \frac{11.3}{T_9^{1/3}} - \frac{2}{3}$$

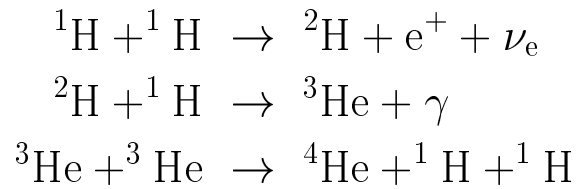
and the lifetime of a proton against destruction is,

$$\tau_p = -\frac{n_p}{dn_p/dt} = \frac{n_p}{2r_{pp}}$$

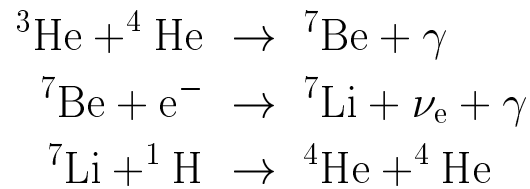
is 6×10^9 yr for Solar central conditions.

To make helium, we need eventually to use 4 protons, two of which need to be converted to neutrons via either positron decays or electron captures. Starting with the fundamental reaction forming a deuteron, there are three main possibilities:

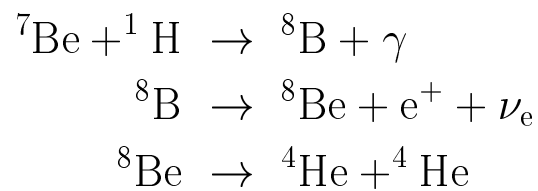
PP-I



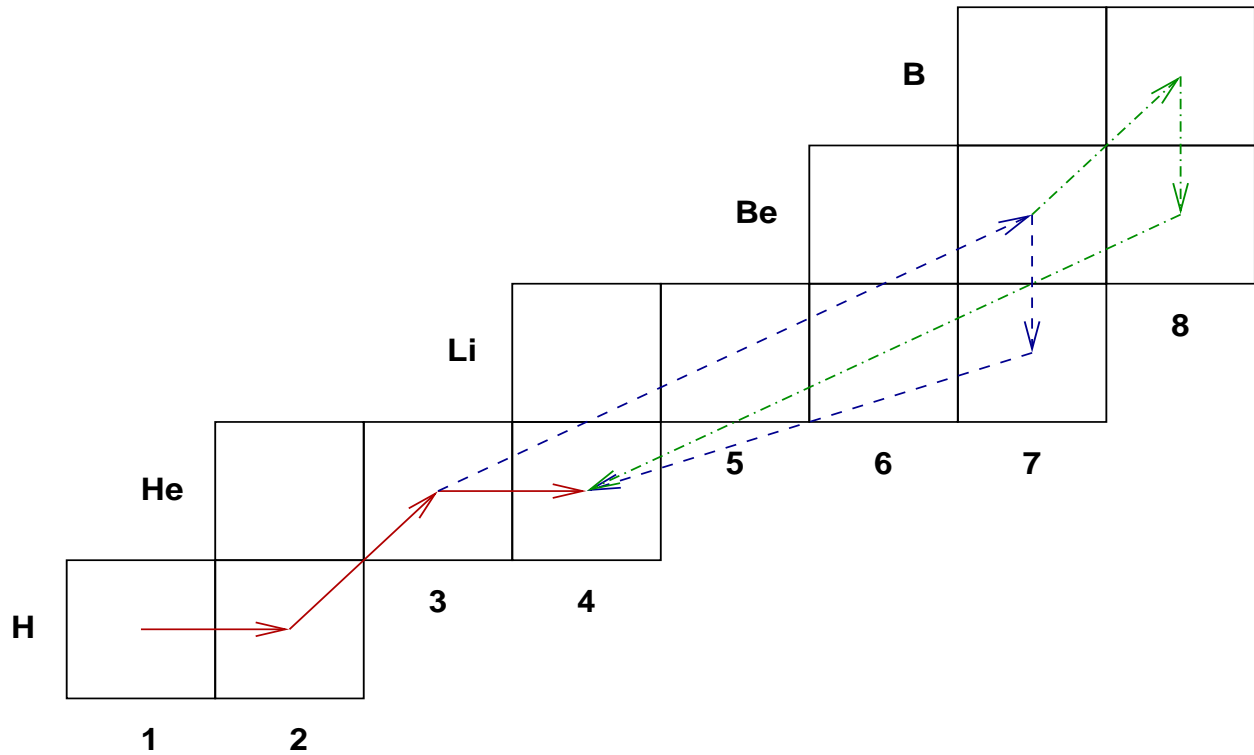
PP-II



PP-III



Graphically,



PP-I is the obvious route, but it requires the somewhat unlikely collision of two relatively short-lived intermediate nuclei. PP-II and PP-III involve splitting a nucleus that is twice as massive as helium.

Relative importance of these chains depends upon the abundance of ${}^3\text{He}$. Termination via PP-I requires two such nuclei to fuse, so the rate of that chain varies as the square of the ${}^3\text{He}$ number density. Find,

- As T increases, equilibrium abundance of ${}^3\text{He}$ decreases.
- With increasing temperature, importance of PP-I compared to PP-II and PP-III decreases.

Although the rate is determined by the slow initial reaction which is common to all three chains, the exact energy yield depends upon the relative importance of PP-I, PP-II and PP-III (because the fraction of the energy lost to ν_e differs). Numerically though, the leading term is,

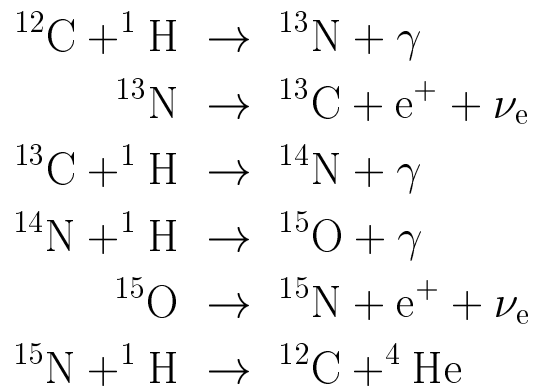
$$\epsilon_{\text{eff}}(\text{pp - chains}) \approx 2.4 \times 10^4 \frac{\rho X^2}{T_9^{2/3}} e^{-3.380/T_9^{1/3}} \text{ erg g}^{-1} \text{ s}^{-1}.$$

Dominant temperature dependence is the exponential arising from $\langle \sigma v \rangle_{pp}$.

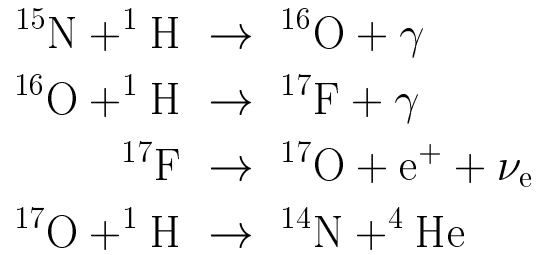
The CNO cycle

The other main route to hydrogen burning is the CNO cycle. This involves carbon, nitrogen and oxygen nuclei in a ‘catalytic’ role. As with the pp reaction, there are several alternative routes.

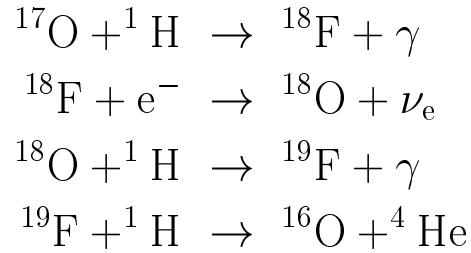
All basically involve a series of proton captures interspersed with positron decays.



Alternative I:



Alternative II:



Description of the CNO cycle is complicated because the timescale to reach equilibrium abundances of the intermediate nuclei is long. Compared to pp, CNO is favoured,

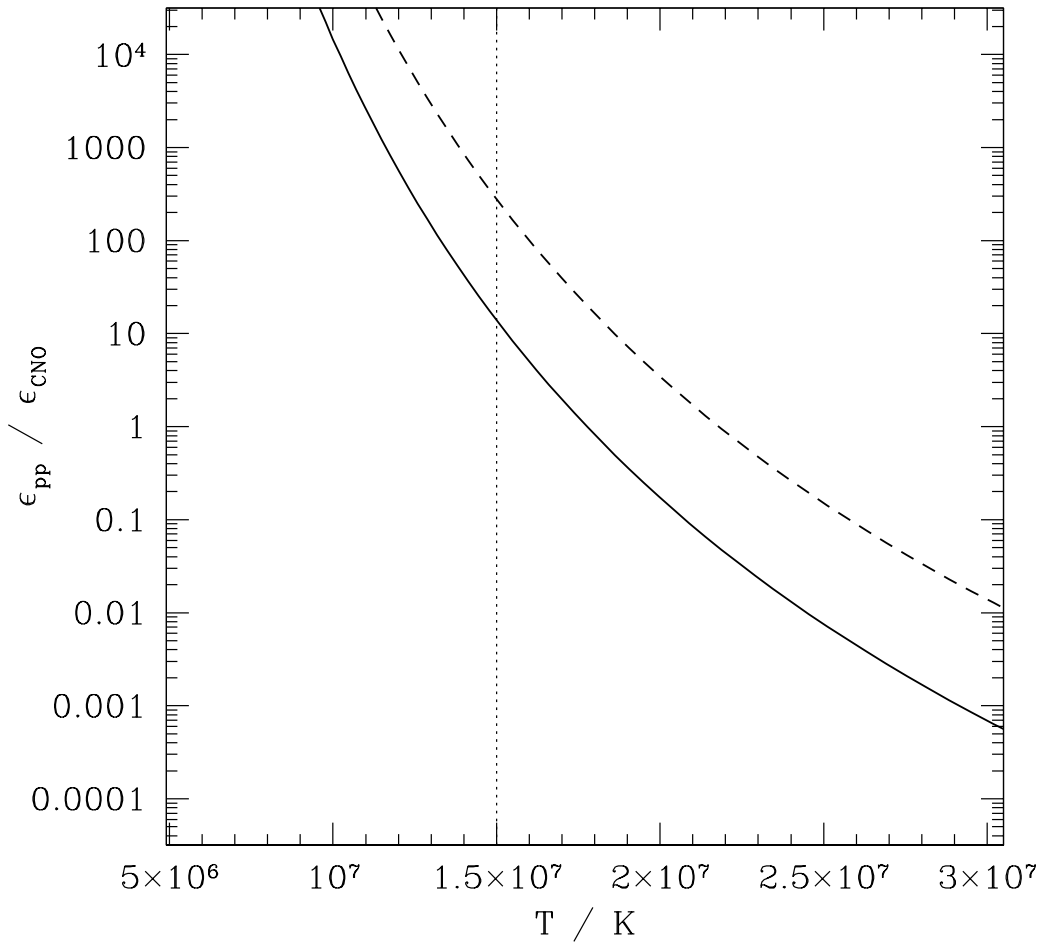
- At higher T , because heavier nuclei are involved.
- At higher metallicity Z .

An approximate rate of energy generation for the CNO cycle is,

$$\epsilon_{\text{CNO}} \approx 4.4 \times 10^{25} \frac{\rho X Z}{T_9^{2/3}} e^{-15.228/T_9^{1/3}} \text{ erg g}^{-1} \text{ s}^{-1}.$$

$\nu \approx 18$ for CNO at $T = 2 \times 10^7$ K.

Relative importance of pp vs CNO cycles:



Solid curve: $X = 0.7, Z = 0.02$. Dashed curve: $X = 0.7, Z = 0.001$.